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## THE BASES OF DIFFERENTIAL GEOMETRY OF VECTOR FIELD IN n-DIMENSIONAL SPACE OF AFFINE CONNECTION.

### Abstract

*Series of invariant straight, hypersurfaces and hyperquadrics for vector field of n-dimensional space of affine connection have been built. The “splitting” of geometric models of vector field in transition from affine space to space of affine connection.*

Classic space of affine connection  $A_{n,n}$  is determined by system  $(n+1)^2$  forms  $\omega^2$  and  $\omega_j^i$ , which ordered to equations of structure

$$\begin{aligned} D\omega^i &= [\omega^k \omega_k^i] + R_{jk}^i [\omega^j \omega^k], \\ D\omega_j^i &= [\omega_i^k \omega_k^j] + R_{ikl}^j [\omega^k \omega^l] \end{aligned} \quad (1)$$

In equations (1) values  $R_{jk}^i$  are sidelong symmetric to upper indexes and in a set they form tensor of steepness of  $A_{n,n}$  space values  $R_{ikl}^j$  are sidelong symmetric to indexes  $k, l$  and they form tensor of steepness  $A_{n,n}$  space.

**Definition.** Vector field in space  $A_{n,n}$  is called the correspondence in which every point  $A(\mathbf{u})$  base of space  $A_{n,n}$  corresponds definite vector  $\vec{v}(u)$  which belongs to n-dimensional affine space  $A_n(\mathbf{u})$  related to moving reper  $T_n = \{\vec{A}(n), e_\alpha(n)\}$ . This space as it is obvious is a layer over point  $A(\mathbf{u})$ .

System of differential equations of vector field in reper of zero order (starting point  $A(\mathbf{u})$  of vector field coincides with the end of vector  $\vec{A}$ , and vector  $\vec{v}$  coincides with  $\vec{v}_n$  en) have the form

$$\omega_n^\alpha = a_{n\beta}^\alpha \omega^\beta \quad (\alpha, \beta, \gamma \dots 1, n) \quad (2)$$

Continuing the system of differential equations (2) we'll receive the system of differential equations of fundamental object of the first order of vector field of space  $A_{n,n}$  in the form

$$d\Lambda_{n\beta}^{\alpha} = \Lambda_{n\gamma}^{\alpha} \omega_{\beta}^{\gamma} - \Lambda_{n\beta}^{\gamma} \omega_{\gamma}^{\alpha} + \Lambda_{n\beta\gamma}^{\alpha} \omega^{\gamma} \quad (3)$$

where

$$\Lambda_{n[\beta\gamma]}^{\alpha} + \Lambda_{n\delta}^{\alpha} R_{\gamma\beta}^{\delta} + R_{n\beta\gamma}^{\alpha} = 0, \quad (4)$$

Continuing the system of differential equations (3) we'll have

$$d\Lambda_{n\beta\gamma}^{\alpha} = \Lambda_{n\delta\gamma}^{\alpha} \omega_{\beta}^{\delta} + \Lambda_{n\beta\delta}^{\gamma} \omega_{\gamma}^{\alpha} - \Lambda_{n\beta\gamma}^{\delta} \omega_{\delta}^{\alpha} + \Lambda_{n\beta\gamma\delta}^{\alpha} \omega^{\delta}. \quad (5)$$

Succession of fundamental objects  $\{\Lambda_{n\beta}^{\alpha}, \Lambda_{n\beta\gamma}^{\alpha}, \Lambda_{n\beta\gamma\delta}^{\alpha}, \dots\}$  lies in the basis of differential geometry of vector field in space  $A_{n,n}$ .

**Remark:** apart from n-dimensional affine space tensors  $\Lambda_{n\beta\gamma}^{\alpha}, \Lambda_{n\beta\gamma\delta}^{\alpha}, \dots$  loose symmetric properties on two down last indexes.

Let's consider values  $\Lambda_{nm}^{\alpha}$ . If their differential equation has form

$$d\Lambda_{nm}^{\alpha} = -\Lambda_{nm}^{\beta} \omega_{\beta}^{\alpha} + \Lambda_{nm}^{\alpha} \omega^{\beta} \quad (6)$$

the straight put by the equation

$$x^{\alpha} = t\Lambda_{nm}^2 \quad (7)$$

relatively to local reper is invariant.

We'll build values  $V_{n\alpha}^{\beta} \Lambda_{n\beta}^{\gamma} = \delta_{\alpha}^{\gamma}$ , in condition, that  $def // \Lambda_{n\beta}^{\alpha} // \neq 0$ .

With the help of value  $V_{n\beta}^{\alpha}$  consider values

$$V^i = -V_{n\beta}^{\alpha} \Lambda_{nm}^{\beta}. \quad (8)$$

Their differential equation has form

$$dV^{\alpha} = -V_{nm}^{\beta} \omega_{\beta}^{\alpha} + V_{\gamma}^{\alpha} \omega^{\gamma} \quad (9)$$

Thus, the straight introduced by the equation

$$x^{\alpha} = tV^{\alpha} \quad (10)$$

is also invariant.

With the help of fundamental objects of the first and second order  $\Lambda_{n\beta}^{\alpha}, \Lambda_{n\beta\gamma}^{\alpha}$  built values in succession

$$\Lambda_{n\beta\alpha}^{\beta} = \Lambda_{n\alpha}^1 \quad d\Lambda_{n\alpha}^1 = \Lambda_{n\beta}^1 \omega_{\alpha}^{\beta} + \Lambda_{n\alpha\beta}^1 \omega^{\beta} \quad (11)$$

$$\Lambda_{n\alpha\beta}^{\beta} = \Lambda_{n\alpha}^2 \quad d\Lambda_{n\alpha}^2 = \Lambda_{n\beta}^2 \omega_{\alpha}^{\beta} + \Lambda_{n\alpha\beta}^2 \omega^{\beta} \quad (12)$$

These values define invariant hypersurface which don't cross the point A:

$$\Lambda_{n\alpha}^1 x^\alpha + 1 = 0 \quad (13)$$

$$\Lambda_{n\alpha}^2 x^\alpha + 1 = 0 \quad (14)$$

In case of space  $A_n$  these hypersurfaces coincide. Thus, we deal with “splitting” of invariant hypersurface which has been mentioned by D.M. Sintsov.

We'll built the following formula

$$\begin{aligned} \Lambda_n &= \Lambda_{n\alpha}^\alpha, \quad \delta \Lambda_n = 0 \\ \Lambda_{n\alpha\beta}^1 &= \Lambda_{n\beta\alpha}^\gamma + \Lambda_{n\gamma}^1 \\ \Lambda_{n\alpha\beta}^2 &= \Lambda_{n\beta\alpha}^\gamma + \Lambda_{n\gamma}^2 \end{aligned} \quad (15)$$

With their help we'll build hyperquadrics

$$\begin{aligned} \Lambda_{n\alpha\beta}^1 x^\alpha x^\beta + 2\Lambda_{n\alpha}^1 x^\alpha + \Lambda &= 0 \\ \Lambda_{n\alpha\beta}^2 x^\alpha x^\beta + 2\Lambda_{n\alpha}^2 x^\alpha + \Lambda &= 0 \end{aligned} \quad (16)$$

In case of n-dimensional affine space these hyperquadrics coincide.