## Linear direct connections in Lie groupoids, curvature and characteristic classes

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## The plan of the talk

1. Linear direct connections  $\tau$  (called also linear quasi-connections) in tangent bundles and in vector bundles. The Teleman's theorem

2. Underlying a usual linear connection  $\nabla^{\tau}$  and a direct proof of this theorem, the curvature of  $\tau$  versus connection of  $\nabla^{\tau}$ .

3. Groupoids point of view and groupoids generalizations.

## 1 Linear direct connections in vector bundles and Teleman's theorem

Nicola Teleman in the papers

N.Teleman, Distance Function, Linear quasi-Connections and Chern Character, June 2004, IHES/M/04/27

N.Teleman, *Direct Connections and Chern Character*, Proceedings of the International Conference in Honor of Jean-Paul Brasselet, Luminy, May 2005,

shows how the Chern character of the tangent bundle of a smooth manifold may be extracted from the geodesic distance function by means of cyclic homology.

The processing has the following steps:

1. Let M be a smooth Riemannian manifold and let

 $r: M \times M \to [0, \infty)$ 

be the induced geodesic distance function.

The function  $r^2$  is smooth on a neighbourhood of the diagonal.

2. Let  $\chi$  be a cut-off smooth monotone decreasing real valued function, identically 1 on a neighbourhood of 0, having support on a sufficiently small interval, so that  $\chi \circ r^2$  be well defined and smooth. For  $x, y \in M$  a linear mapping

$$A(y,x):T_xM\to T_yM$$

is given by the formula

$$A(y,x)\left(\sum_{i}\xi^{i}\frac{\partial}{\partial x^{i}}\right) = \sum_{i,j,k}\xi^{i}\frac{\partial^{2}\left(\chi\circ r^{2}\right)\left(x,y\right)}{\partial x^{i}\partial y^{j}}g^{jk}\left(y\right)\frac{\partial}{\partial y^{k}}$$

(A(y, x) is independent of the local coordinates).

For sufficiently close points x, y,

-A(y, x) is an isomorphism and

-A(x,x) is the identity.

Therefore A is a linear direct connection (=linear quasi-connection), with respect to the definition below.

3. With the object A there is associated the function  $\Phi_k : U_{k+1} \to \mathbb{R}$ , where  $U_{k+1}$  is a neighbourhood of the diagonal in  $M^{k+1}$ 

 $\Phi_{k}(x_{0}, x_{1}, ..., x_{k}) := Trace \ A(x_{0}, x_{1}) \circ A(x_{1}, x_{2}) \circ ... \circ A(x_{k-1}, x_{k}) \circ A(x_{k}, x_{0}).$ 

4. Next, N.Teleman studies the function  $\Phi_k$  in the context of cyclic homology:

— firstly, he notices that  $\Phi_k$ , k =even, is a cyclic cycle over the algebra  $\mathcal{A} = C^{\infty}(M)$ ,

— secondly, he uses – the Connes' isomorphism which associates with  $\Phi_k$  a closed differential form

$$\Omega\left(\Phi_{k}\right)\left(x\right) = \frac{1}{k!} \sum_{i_{1},i_{2},\dots,i_{k}} \frac{\partial}{\partial x_{1}^{i_{1}}} \frac{\partial}{\partial x_{2}^{i_{2}}} \dots \frac{\partial}{\partial x_{k}^{i_{k}}} \Phi_{k}\left(x_{0},x_{1},\dots,x_{k}\right)_{x_{0}=x_{1}=\dots=x_{k}=x} dx^{i_{1}} \wedge \dots \wedge dx^{i_{k}},$$

(we use the same local coordinate system on each factor).

— thirdly, he proves

**Theorem 1** The top degree component of the cyclic homology class of  $\Phi_k$  is equal to

$$\left[\Omega\left(\Phi_{2k}\right)\right] = c \cdot Ch_k\left(M\right)$$

where c is a constant and  $Ch_k(M)$  is the k-component of the Chern character of the tangent bundle of M.

The object A is a particular case of the linear direct connection introduced by N.Teleman.

**Definition 2** Let E be a real or complex smooth vector bundle over the manifold M. A **linear direct connection**  $\tau$  in E consists of assigning to any two points  $x, y \in M$ , sufficiently close one to each other, an isomorphism

$$\tau\left(y,x\right):E_{|x}\to E_{|y},$$

such that

$$\tau\left(x,x\right) = id,$$

and  $\tau(y, x)$  depends smoothly on the pair x, y.

The parallel transport defined by a usual linear connection in E along the small geodesics of an affine connection in M induces a linear direct connection in E (see for example A.Connes and H.Moscovici, "*Cyclic cohomology, the Novikov conjecture and hyperbolic groups*", Topology 29, n 3 345-388, 1990).

-i) As for A with  $\tau$  there is associated the function  $\Phi_k$  by the formula

 $\Phi_k(x_0, x_1, ..., x_k) := Trace \ \tau \ (x_0, x_1) \circ \tau \ (x_1, x_2) \circ ... \circ \tau \ (x_{k-1}, x_k) \circ \tau \ (x_k, x_0) \,.$ 

The function

 $\Phi_2(x_0, x_1, x_2) = Trace \ \tau(x_0, x_1) \circ \tau(x_1, x_2) \circ \tau(x_2, x_0)$ 

plays a role of the *curvature* of  $\tau$  and the differential form  $\Omega(\Phi_2)$  - the *curvature* form of  $\tau$ .

-ii) Any two smooth linear direct connections in a smooth vector bundle are smoothly homotopic. The results above implay

**Theorem 3** (N. Teleman) For any smooth linear direct connection  $\tau$  in the smooth vector bundle E over the manifold M,

- -i)  $\Phi_k$ , k = even, is a cyclic cycle over the algebra  $C^{\infty}(M)$ ,
- -*ii*) the cohomology class of  $\Omega(\Phi_{2k})$  is (up to a multiplicative constant) is the k-component of the Chern character of E.

# 2 Underlying linear connection $\nabla^{\tau}$ and a direct proof of this theorem

In the paper

J.Kubarski, N.Teleman, *Linear direct connections*, Banach Center Publications, 2007, in print,

we study the geometry of direct connections  $\tau$ :

• we construct the "infinitesimal part"  $\nabla^{\tau}$  and show that  $\nabla^{\tau}$  is a usual linear connection. We next determine the curvature tensor R of  $\nabla^{\tau}$  and show that the **equality of differential forms** holds

$$\Omega\left(\Phi_{2k}\right) = c \cdot Tr \ R^k.$$

We intend to extract from a direct connection its infinitesimal part along the diagonal.

**Definition 4** Let X be a smooth tangent field over M and  $\phi$  a smooth section in E. Let  $x_0$  be an arbitrary point in M and let  $\gamma : (-\varepsilon, \varepsilon) \longrightarrow M$  be an integral path of the field X with the initial condition  $\gamma(0) = x_0$ .

We define

$$\nabla_{X(x_0)}^{\tau}(\phi) = \frac{d}{dt} \left\{ \tau(\gamma(0), \gamma(t)) \left( \phi(\gamma(t)) \right) \right\}_{|t=0} \in E_{|x_0}.$$

**Theorem 5** The right hand side of the above formula depends only on the value of X at  $x_0$ . The operator  $\nabla^{\tau}_{X(x_0)}(\phi)$  is a usual linear connection in E.

We intend to describe  $\nabla^{\tau}_{X(x_0)}(\phi)$  locally.

Let  $(x^1, x^2, ..., x^m)$  (dim M = m) be a local coordinate system on an open neighborhood  $\mathcal{V}$  of a point  $x_0$ . Using the same local coordinate system on both factors of the direct product  $M \times M$ , any point  $(x, y) \in \mathcal{V} \times \mathcal{V}$  will be given by local coordinates  $(x^1, x^2, ..., x^m | y^1, y^2, ..., y^m)$ .

**Theorem 6** Let  $\{e_1, e_2, ..., e_n\}$  be a local frame in E over  $\mathcal{V}$ . Let  $\tau(x|y)$  be the matrix describing locally the direct connection  $\tau$ :

$$\tau(x|y) = \|\tau_i^j(x|y)\| \in M_{n,n}(\mathbb{K}),$$
  
$$\tau(x,y) \left(e_i(y)\right) = \sum_j \tau_i^j(x|y) \cdot e_j(x), \quad \tau_i^j(x|x) = \delta_i^j.$$

Then the coefficients  $\Gamma^j_{i,\alpha}$  of the connection  $\nabla^{\tau}$  are given locally by

$$\nabla_{\frac{\partial}{\partial x^{\alpha}}}^{\tau} e_i = \sum_j \Gamma_{i,\alpha}^j e_j,$$

where

$$\Gamma^j_{i,\alpha}(x) = \frac{\partial}{\partial y^\alpha} \tau^j_i(x^1, x^2, ..., x^m | y^1, y^2, ..., y^m)_{y=x}$$

In conclusion, representing the tangent field X locally

$$X(x) = \sum_{\alpha} X^{\alpha}(x) \cdot \frac{\partial}{\partial x^{\alpha}},$$

one has the formula

$$\nabla_{X(x_0)}^{\tau}(\sum_i \phi^i e_i) = \sum_{\alpha=1}^m \{ \sum_{i,j} \Gamma_{i,\alpha}^j(x_0) \cdot X^{\alpha}(x_0) \cdot \phi^i(x_0) e_j(x_0) \} + \sum_i (d\phi^i)(X)(x_0) e_i(x_0).$$

**Remark 7** The above formula also show that  $\nabla^{\tau}$  is a linear connection in the vector bundle E. The linear connection  $\nabla^{\tau}$  will be called associated, or underlying, linear connection to the direct connection  $\tau$ .

**Proposition 8** Let  $R = (\nabla^{\tau})^2$  be the curvature tensor of the connection  $\nabla^{\tau}$ . The components of the curvature R are

$$\begin{split} R^{j}_{i\alpha\beta}(x) &= \frac{\partial}{\partial x^{\alpha}} \Gamma^{j}_{i\beta}(x) - \frac{\partial}{\partial x^{\beta}} \Gamma^{j}_{i\alpha}(x) + \Gamma^{j}_{k\alpha}(x) \cdot \Gamma^{k}_{i\beta}(x) - \Gamma^{j}_{k\beta}(x) \cdot \Gamma^{k}_{i\alpha}(x) \\ &= \frac{\partial^{2}}{\partial x^{\alpha} \partial y^{\beta}} \tau^{j}_{i}(x|y)_{y=x} - \frac{\partial^{2}}{\partial x^{\beta} \partial y^{\alpha}} \tau^{j}_{i}(x|y)_{y=x} + \\ &+ \frac{\partial}{\partial y^{\alpha}} \tau^{j}_{k}(x|y)_{y=x} \cdot \frac{\partial}{\partial y^{\beta}} \tau^{k}_{i}(x|y)_{y=x} - \frac{\partial}{\partial y^{\beta}} \tau^{j}_{k}(x|y)_{y=x} \cdot \frac{\partial}{\partial y^{\alpha}} \tau^{k}_{i}(x|y)_{y=x}. \end{split}$$

**Corollary 9** The curvature form R of the underlying linear connection  $\nabla^{\tau}$ , associated to the direct connection  $\tau$ , is given by

$$\begin{split} R &= (\frac{\partial^2}{\partial x^{\alpha} \partial y^{\beta}} \tau_i^j(x|y)_{y=x} - \frac{\partial^2}{\partial x^{\beta} \partial y^{\alpha}} \tau_i^j(x|y)_{y=x} + \\ &+ \frac{\partial}{\partial y^{\alpha}} \tau_k^j(x|y)_{y=x} \cdot \frac{\partial}{\partial y^{\beta}} \tau_i^k(x|y)_{y=x} - \frac{\partial}{\partial y^{\beta}} \tau_k^j(x|y)_{y=x} \cdot \frac{\partial}{\partial y^{\alpha}} \tau_i^k(x|y)_{y=x}) dx^{\alpha} \wedge dx^{\beta}. \end{split}$$

Although,  $\tau(x, y) = (\tau(y, x))^{-1}$  is not true in general, it is true, however, that it holds infinitesimally. In fact, we have the

**Proposition 10** For any direct connection  $\tau$ , its matrix components satisfy the identities

-i)

$$\frac{\partial}{\partial x^{\alpha}}\tau_i^j(x|y)_{y=x} + \frac{\partial}{\partial y^{\alpha}}\tau_i^j(x|y)_{y=x} = 0.$$

-ii)

$$\frac{\partial}{\partial x^{\alpha}} \{\tau(x|y) \circ \tau(y|x)\}_{y=x} = 0 = \frac{\partial}{\partial y^{\alpha}} \{\tau(x|y) \circ \tau(y|x)\}_{y=x}.$$

As  $\tau(x|x) = Id$ , we get that the directional derivative  $\left(\frac{\partial}{\partial x^{\alpha}} + \frac{\partial}{\partial y^{\alpha}}\right)$  of  $\tau$  along the diagonal **vanishes**. This proves -i). The second identity is a consequence of the first.

The above properties of any direct connection are fundamental for comparing the curvature tensor R to the differential form  $\Omega(\Phi_{2k}^{\tau})$ .

We obtain an important explicit link between  $\Omega(\Phi_{2k})$  and the classical Chern-Weil forms, at the level of **differential forms** rather than **cohomology** classes.

**Theorem 11** Let  $\tau$  be a direct connection and let  $\nabla^{\tau}$  be its underlying linear connection. Then

$$\Omega(\Phi_2^{\tau}) = \frac{1}{4} \cdot Tr \ R,$$

and more generally,

$$\Omega(\Phi_{2k}^{\tau}) = \frac{1}{(2k)!} \cdot \frac{1}{2^k} \cdot Tr \ R^k,$$

where  $R = (\nabla^{\tau})^2$  is the curvature of the underlying linear connection  $\nabla^{\tau}$ .

In consequence, the mentioned above Teleman's theorem follows from this directly.

### 3 Groupoids point of view and groupoids generalizations

N.Teleman in yours papers said:

"The arguments discussed here may be extended to the language of groupoids". My further talk is the first step in this direction.

#### **3.1** Direct connections and the Lie groupoid GL(E)

Let E be a real or complex smooth vector bundle over the manifold M. Consider the **transitive Lie groupoid** 

$$\Phi = GL\left(E\right)$$

of all linear fibre isomorphisms  $h: E_{|x} \to E_{|y}$  of the vector bundle E, with the source  $\alpha$ ,  $\alpha(h) = x$ , and the target  $\beta$ ,  $\beta(h) = y$ , and the unit  $u_y = \mathrm{id}_{E_{|y}}$ . The mappings

$$\alpha, \beta: \Phi \to M, \quad (\alpha, \beta): \Phi \to M \times M$$

are submersions, the injection

$$u: M \to \Phi, y \to u_y,$$

is smooth, and the partial multiplication

$$\cdot: \Phi \times_{(a,\beta)} \Phi \to \Phi, \ (g,h) \longmapsto gh,$$

is also smooth. and GL(E) be a Lie groupoid of linear fibre isomorphisms.

**Remark 12** A linear direct connection in a vector bundle E is equivalently a smooth mapping

$$\tau: U \to GL(E)$$

where  $U \subset M \times M$  is an open neighborhood of the diagonal  $\Delta = \{(x, x); x \in M\}$ , such that

$$\tau\left(x,y\right):E_{|y}\to E_{|x}$$

i.e.

$$\alpha \circ \tau (x, y) = y, \quad \beta \circ \tau (y, x) = x,$$

and

$$\tau(x,x) = \mathrm{id}: E_{|x} \to E_{|x}.$$

#### 3.2 Lie Groupoids and point of view of linear direct connections and the using of the Lie algebroids

According to the Pradines definition, the **Lie algebroid** of an arbitrary transitive Lie groupoid  $\Phi$  is equal to the vector bundle

$$A\left(\Phi\right) = u^*\left(T^{\alpha}\Phi\right)$$

where  $u: M \to \Phi, y \to u_y$ , and  $T^{\alpha}\Phi = \ker \alpha_*$ , equipped with the suitable structures: the bracket of cross-sections  $[\![\xi, \eta]\!], \xi, \eta \in SecA(\Phi)$  is defined in the following way. The cross-sections  $\xi, \eta$  can be extended to right invariant vector fields  $\xi', \eta'$  on  $\Phi$ , their usual bracket  $[\xi', \eta']$  is invariant too, so it determines a cross-section of  $u^*(T^{\alpha}\Phi)$  denoting by  $[\![\xi, \eta]\!]$ . The anchor is defined as the restriction of  $\beta_*$ .

We recall the definition of a Lie algebroid.

**Definition 13** By a Lie algebroid on a manifold M we mean a system

$$A = (A, \llbracket \cdot, \cdot \rrbracket, \gamma) \tag{1}$$

consisting of a vector bundle A (over M) and mappings

$$\llbracket \cdot, \cdot \rrbracket : \operatorname{Sec} A \times \operatorname{Sec} A \longrightarrow \operatorname{Sec} A, \quad \gamma : A \longrightarrow TM,$$

such that

- (i)  $(\text{Sec } A, \llbracket \cdot, \cdot \rrbracket)$  is an  $\mathbb{R}$ -Lie algebra,
- (ii)  $\gamma$ , called an A nchor, is a homomorphism of vector bundles,
- (*iii*) Sec  $\gamma$  : Sec  $\gamma \longrightarrow \mathfrak{X}(M)$ ,  $\xi \longmapsto \gamma \circ \xi$ , is a homomorphism of Lie algebras,
- $(iv) \ \llbracket \xi, f \cdot \eta \rrbracket = f \cdot \llbracket \xi, \eta \rrbracket + (\gamma \circ \xi) \left( f \right) \cdot \eta \text{ for } f \in C^{\infty} \left( M \right), \, \xi, \eta \in \operatorname{Sec} A.$

Lie algebroid (1) is called *transitive* if  $\gamma$  is an epimorphism.  $\boldsymbol{g} = \ker \gamma$  is a vector bundle, called the adjoint of (1), and the short exact sequence

$$0 \longrightarrow \boldsymbol{g} \hookrightarrow A \xrightarrow{\gamma} E \longrightarrow 0 \tag{2}$$

is called the Atiyah sequence of (1).

**Example 14** The following are simple fundamental examples of transitive Lie algebroids:

- $(1^0)$  Finitely dimensional Lie algebra.
- $(2^0)$  Tangent bundle TM to a manifold M with the bracket  $[\cdot, \cdot]$  of vector fields and  $id_{TM}$  as an anchor.
- (3<sup>0</sup>) Trivial Lie algebroid  $TM \times \mathfrak{g}$  (Ngo-Van-Que) where  $\mathfrak{g}$  is as in (1<sup>o</sup>). The bracket is defined by the formula,

 $\llbracket (X,\sigma), (Y,\eta) \rrbracket = ([X,Y], \mathcal{L}_X \eta - \mathcal{L}_Y \sigma + [\sigma,\eta]),$ 

 $X, Y \in \mathcal{X}(M), \ \sigma, \eta: M \to \mathfrak{g}$ , and the anchor is the projection  $TM \times \mathfrak{g} \to TM$ .

 $(4^0)$  Bundle of jets  $J^kTM$  (P.Libermann).

(5<sup>0</sup>) General form (K.Mackenzie, J.Kubarski). Let a system  $(\boldsymbol{g}, \nabla, \Omega_b)$  be given, consisting of a Lie algebra bundle  $\boldsymbol{g}$  on a manifold M, a covariant derivative  $\nabla$  in  $\boldsymbol{g}$  and a 2-form  $\Omega_b \in \Omega^2(M, \boldsymbol{g})$  on M with values in  $\boldsymbol{g}$ , fulfilling the conditions:

(i) 
$$\nabla^2 \sigma = -[\Omega_b, \sigma], \quad \sigma \in \operatorname{Sec} \boldsymbol{g},$$
  
(ii)  $\nabla_X[\sigma, \eta] = [\nabla_X \sigma, \eta] + [\sigma, \nabla_X \eta], \quad X \in X(M), \quad \sigma, \eta \in \operatorname{Sec} \boldsymbol{g},$   
(iii)  $\nabla \Omega_b = 0.$ 

Then  $TM \oplus \boldsymbol{g}$  forms a transitive Lie algebroid with the bracket defined by

 $\llbracket (X,\sigma), (Y,\eta) \rrbracket = ([X,Y], -\Omega_b(X,Y) + \nabla_X \eta - \nabla_Y \sigma + [\sigma,\eta]),$ 

the anchor being the projection onto the first component.

Every transitive Lie algebroid is — up to an isomorphism — of this form.

**Example 15** The following are important examples of transitive Lie algebroids:

- (6<sup>0</sup>) The Lie algebroid A(P) = TP/G of a G-principal bundle P (K.Mackenzie, J.Kubarski).
- (7<sup>0</sup>) The Lie algebroid CDO(E) of covariant differential operators on a vector bundle E (K.Mackenzie). Another isomorphic construction of this object is the Lie algebroid A(E) of a vector bundle E (J.Kubarski), here the fibre  $A(E)_{|x}$  is the space of linear homomorphisms  $l : \text{Sec } E \to E_{|x}$  such that there exists a vector  $u \in T_x M$  for which  $l(f \cdot \nu) = f(x) \cdot l(\nu) + u(f) \cdot \nu(x)$ ,  $f \in C^{\infty}(M), \nu \in \text{Sec}(E)$ .
- ( $\delta^{0}$ ) The Lie algebroid  $A(\Phi) := i^{*}T^{\alpha}\Phi$  of a Lie groupoid  $\Phi$  (J.Pradines). (Remark: if  $\Phi = GL(E)$  is the Lie algebroid of a all linear fibre isomorphisms of fibres of E then  $A(E) = A(\Phi)$ ).
- (9) The Lie algebroid  $A(M, \mathcal{F})$  of a transversally complete foliation  $(M, \mathcal{F})$ (P.Molino); in particular,
- $(10^{0})$  the Lie algebroid A(G; H) of the foliation of left cosets of a Lie group G by a nonclosed connected Lie subgroup  $H \subset G$  (for the construction independent of the theory of transversally complete foliations, see J.Kubarski).

There are many sources of nontransitive Lie algebroids: Lie equations, Differential groupoids, Poisson manifolds, etc.

Let  $\Phi = GL(E)$  be the Lie groupoid of all linear fibre isomorphisms of fibres of E.

For  $y \in M$  the submanifold  $\Phi_y = GL(E)_y \subset GL(E)$  of all elements  $u \in GL(E)$  for which  $\alpha(u) = y$ ,

$$GL(E)_{y} = \alpha^{-1}(y),$$

is a  $GL(E_y)$ -principal fibre bundle.

- Lie algebroid of the Lie groupoid is the infinitesimal object and play analogous role to that of Lie algebras for Lie groups.
- The space [Lie algebra] of global cross-sections  $Sec(A(\Phi))$ ,  $\Phi = GL(E)$ where E is a vector bundle, is naturally isomorphic to the Lie algebra of all Covariant Derivative Operators, i.e. to the space of differential operators of the rank  $\leq 1$

$$\mathfrak{L}:SecE\to SecE$$

such that  $\mathfrak{L}(f \cdot \xi) = f \cdot \mathfrak{L}(\xi) + X(f) \cdot \xi$ , for a vector field X called the anchor of  $\mathfrak{L}, f \in C^{\infty}(M), \xi \in SecE$ .

Let  $\tau : (M \times M)_{|U} \to GL(E)$  (where  $U \subset M \times M$  is an open neighbourhood of the diagonal  $\Delta = \{(x, x); x \in M\}$ ) be a linear quasi-connection,

$$\tau_{(x,y)}: E_{|y} \to E_{|x|}$$

so  $\alpha(\tau_{(x,y)}) = y$  and  $\beta(\tau_{(x,y)}) = x$  and let  $\nabla^{\tau}$  be the underlying linear connection of  $\tau$  in E.

Now, we fix y and take

$$\tau(\cdot, y): M \to GL(E)_y, \quad x \longmapsto \tau(x, y).$$

It is a smooth mapping – such that  $\beta \circ \tau (\cdot, y) = id$ . Therefore the composition of the differential

$$\tau\left(\cdot,y\right)_{*x}:T_{x}M\to T_{\tau\left(x,y\right)}\left(GL\left(E\right)_{y}\right)$$

with the differential of  $\beta | GL(E)_y \to M$  is identity

$$\mathrm{id}: T_x M \xrightarrow{\tau(\cdot,y)_{*x}} T_{\tau(x,y)} \left( GL\left(E\right)_y \right) \xrightarrow{\beta_*} T_x M.$$

Taking x = y and using the fact  $\tau(y, y) = u_y = \mathrm{id}_{E_y}$  we see that

$$\tau\left(\cdot,y\right)_{*y}:T_{y}M\to T_{u_{y}}\left(GL\left(E\right)_{y}\right).$$

Therefore  $\tau$  determines a usual connection

$$\bar{\nabla}^{\tau} : TM \to u^* (T^{\alpha} \Phi)$$
$$\bar{\nabla}^{\tau} (v_y) = \tau (\cdot, y)_{*y} (v_y).$$

in the Lie algebroid  $u^{*}(T^{\alpha}\Phi)$  ( $\Phi = GL(E)$ ), i.e. a splitting of the Atiyah sequence

$$0 \to \boldsymbol{g} \to A(\Phi) \xrightarrow[\nabla^{\tau}]{\beta_*} TM \to 0.$$

 $\overline{\nabla}^{\tau}$  is the "usual covariant derivative" since the anchor of the Covariant Derivative Operator  $\overline{\nabla}^{\tau}(X) : SecE \to SecE$  is just equal to X, therefore noticing  $\overline{\nabla}^{\tau}(X)(\xi)$  in the form

 $\bar{\nabla}_{X}^{\tau}\left(\xi\right)$ 

the usual axioms for covariant derivative are fulfilled.

**Theorem 16**  $\overline{\nabla}^{\tau} = \nabla^{\tau}$ , *i.e. the connection*  $\overline{\nabla}^{\tau}$  *is equal to the underlying linear connection of*  $\tau$  *in* E.

**Proof.** (a sketch) Since we need to prove it at any point  $y \in M$  so we can prove it locally for  $E = M \times \mathbb{R}^n$  and  $M = \mathbb{R}^m$ . Then  $GL(E) = M \times GL(\mathbb{R}^n) \times M$ ,  $\alpha^{-1}(y) = GL(E)_y = M \times GL(\mathbb{R}^n) \times \{y\}$ . Let  $\{e_i\}_{i=1}^n$  be a trivial local basis of E, then the induced linear connection  $\nabla^{\tau}$  is determined by

$$\nabla_{\frac{\partial}{\partial x^{k}}|y}^{\tau}e_{i} = \frac{\partial \tau_{i}^{j}}{\partial x^{m+k}}(y,y) \cdot e_{j}.$$

We can obtain the same results for  $\overline{\nabla}^{\tau}$ .

#### 3.3 Groupoids generalization

The above consideration has "groupoids sense" so we can it generalize to any transitive Lie groupoids.

Let  $\Phi$  be an arbitrary transitive Lie groupoid with the anchor  $\alpha$  and the target  $\beta$ . We denote by  $u_y$  the unit of  $\Phi$  at y.

**Definition 17** By a linear direct connection in  $\Phi$  we mean a mapping

$$\tau: (M \times M)_{|U} \to \Phi,$$

such that

$$\alpha \circ \tau (x, y) = y, \quad \beta \circ \tau (x, y) = x,$$

and

$$\tau\left(x,x\right) = u_x.$$

For y the submanifold  $\Phi_y \subset \Phi$  of all elements  $h \in \Phi$  for which  $\alpha(h) = y$  $(\Phi_y = \alpha^{-1}(y))$  is a  $\Phi_y^y$ -principal fibre bundle where

$$\Phi_{u}^{y} = \{h \in \Phi; \ \alpha(h) = \beta(h) = y\}$$

is the isotropy Lie algebra of  $\Phi$  at y. Now, we fix y and take

$$\tau(\cdot, y): M \to \Phi_y, \quad x \longmapsto \tau(x, y).$$

It is a smooth mapping such that

$$\beta \circ \tau (\cdot, y) = \mathrm{id}$$
.

Taking the differential

$$\tau\left(\cdot,y\right)_{*x}:T_xM\to T_{\tau(x,y)}\left(\Phi_y\right)$$

such that the composition with the differential of  $\beta | \Phi_y \to M$  is identity

$$\mathrm{id}: T_x M \xrightarrow{\tau(\cdot,y)_{*x}} T_{\tau(x,y)} \left( \Phi_y \right) \xrightarrow{\beta_*} T_x M$$

and taking x = y and using the fact  $\tau(y, y) = u_y$  we see that

$$\tau\left(\cdot,y\right)_{*y}:T_{y}M\to T_{u_{y}}\left(\Phi_{y}\right)=A\left(\Phi\right)_{|y},$$

where  $A(\Phi)$  is the Lie algebroid of the Lie groupoid  $\Phi$ .

Therefore  $\tau$  determines a splitting of the Atiyah sequence of  $\Phi$ 

$$0 \to \mathbf{g} \to A\left(\Phi\right) \xrightarrow[]{\beta_*}{\overleftarrow{\nabla^{\tau}}} TM \to 0,$$

i.e. a usual connection in the Lie algebroid  $A(\Phi) = u^*(T^{\alpha}\Phi)$ ,

$$\nabla^{\tau} : TM \to u^* \left( T^{\alpha} \Phi \right) = A \left( \Phi \right)$$
$$\nabla^{\tau} \left( v_y \right) = \tau \left( \cdot, y \right)_{*y} \left( v_y \right)$$

The connection  $\nabla^{\tau}$  will be called the **underlying linear connection of the linear direct connection**  $\tau$ .

Now we can ask on a very important question:

• How can we reconstruct the curvature tensor of  $\nabla$  from the linear direct connection in the Lie groupoid  $\Phi$ ? And next how can we reconstruct the Chern-Weil homomorphism of Lie groupoids  $\Phi$  (i.e. equivalently of the principal bundle  $\Phi_y$ ) from arbitrary taken linear direct connection  $\tau$ ?

## 3.4 Curvature tensor of the linear direct connection in transitive Lie groupoids

Take any transitive Lie groupoid  $\Phi$  and its Lie algebroid  $A(\Phi)$  with the Atiyah sequence

$$0 \rightarrow \boldsymbol{g} \rightarrow A(\Phi) \longrightarrow TM \rightarrow 0.$$

The fibre of  $\boldsymbol{g}$  at x

$$\boldsymbol{g}_{|x} = T_{u_x} \Phi_x^x$$

is the **right** Lie algebra of the structural Lie group  $\Phi_x^x$ . For a linear direct connection  $\tau$  in  $\Phi$  denote by

$$\nabla^{\tau}: TM \to A(\Phi)$$

the underlying linear connection in the Lie algebroid  $A(\Phi)$  induced by  $\tau$ . Consider the curvature tensor  $\Omega^{\tau} \in \Omega^2(M; \boldsymbol{g})$  of  $\nabla^{\tau}$ 

$$\Omega^{\tau}(X,Y) = \llbracket \nabla_X^{\tau}, \nabla_Y^{\tau} \rrbracket - \nabla_{[X,Y]}^{\tau}.$$

The linear direct connection  $\tau$  determines the mapping

$$\Psi_k^{\tau} : \left(\underbrace{M \times \dots \times M}_{k+1}\right)_{|U} \to \Phi,$$
  
$$\Psi_k^{\tau}(x_0, x_1, \dots, x_k) = \tau(x_0, x_1) \cdot \tau(x_1, x_2) \cdot \dots \cdot \tau(x_{k-1}, x_k) \cdot \tau(x_k, x_0)$$

having the values in the associated Lie group bundle,

$$\Psi_k^{\tau}(x_0, x_1, ..., x_k) \in \Phi_{x_0}^{x_0}.$$

For example, for k = 2, the function

$$\Psi_2^{\tau} : (M \times M \times M)_{|U} \to \Phi,$$
  
$$\Psi_2^{\tau} (x_0, x_1, x_2) = \tau (x_0, x_1) \cdot \tau (x_1, x_2) \cdot \tau (x_2, x_0)$$

is called the *curvature* of  $\tau$ .

Analogously to the previous cases we can associate some **differential form** to the function  $\Psi_k$ . Namely, fixing a point  $x_0$  we define

$$\Psi_{k}^{\tau}(x_{0}):\left(\underbrace{M\times\ldots\times M}_{k}\right)_{|U}\to\Phi_{x_{0}}^{x_{0}},$$
$$(x_{1},\ldots,x_{k})\longmapsto\Psi_{k}^{\tau}(x_{0},x_{1},\ldots,x_{k})\in\Phi_{x_{0}}^{x_{0}}.$$

Next, we take a coordinate system  $(x^1, ..., x^m)$  (dim M = m) on an open neighborhood  $\mathcal{V}$  of the point  $x_0$ . Using the same local coordinate system on each factors of the direct product  $M \times \cdots \times M$  we take for  $(x_1, ..., x_k) \in \mathcal{V} \times ... \times \mathcal{V}$ 

$$\frac{\partial}{\partial x_k^j} \Psi_k^\tau(x_0) \in T_{\Psi_k(x_0, x_1, \dots, x_k)} \Phi_{x_0}^{x_0}.$$

This vector we can translate via **right** translation to the unit. Let  $r_h : \Phi_{x_0}^{x_0} \to \Phi_{x_0}^{x_0}$  denote the right translation on the element h,  $r_h(z) = z \cdot h$ .

$$\frac{\partial}{\partial \tilde{x}_k^j} \Psi_k^\tau(x_0) := \left( r_{(\Psi_k(x_0, x_1, \dots, x_k))^{-1}} \right)_{*\Psi_k(x_0, x_1, \dots, x_k)} \left( \frac{\partial}{\partial x_k^j} \Psi_k^\tau(x_0) \right) \in T_{u_{x_0}}\left( \Phi_{x_0}^{x_0} \right) = \boldsymbol{g}_{|x_0}.$$

The function obtained

$$(x_1, ..., x_{k-1}, x_k) \longmapsto \frac{\partial}{\partial \tilde{x}_k^j} \Psi_k^{\tau}(x_0) \in \boldsymbol{g}_{|x_0|}$$

can be differentiated usually as a vector valued function.

$$(x_1, ..., x_{k-1}, x_k) \longmapsto \frac{\partial}{\partial x_1^{i_1}} \frac{\partial}{\partial x_2^{i_2}} ... \frac{\partial}{\partial x_{k-1}^{i_{k-1}}} \frac{\partial}{\partial \tilde{x}_k^{i_k}} \Psi_k^{\tau} (x_0, x_1, ..., x_k) \in \boldsymbol{g}_{|x_0}.$$

We put

$$\Omega\left(\Psi_{k}^{\tau}\right)\left(x\right) = \frac{1}{k!} \sum_{i_{1},i_{2},\dots,i_{k}} \frac{\partial}{\partial x_{1}^{i_{1}}} \frac{\partial}{\partial x_{2}^{i_{2}}} \frac{\partial}{\partial \tilde{x}_{k}^{i_{k}}} \Psi_{k}^{\tau}\left(x_{0},x_{1},\dots,x_{k}\right)_{x_{0}=x_{1}=\dots=x_{k}=x} dx^{i_{1}} \wedge \dots \wedge dx^{i_{k}}.$$

It is a k-form on M with values in the vector bundle  $\boldsymbol{g}$ 

$$\Omega\left(\Psi_{k}^{\tau}\right)\in\Omega^{k}\left(M;\boldsymbol{g}\right)$$

Considering k = 2 we obtain a 2-form with values in  $\boldsymbol{g}$ ,

$$\Omega\left(\Psi_{2}^{\tau}\right)\in\Omega^{2}\left(M;\boldsymbol{g}\right),$$

called the **curvature form** of  $\tau$ .

The fundamental role is playing by the following

**Theorem 18** For an arbitrary linear direct connection  $\tau : (M \times M)_{|U} \to \Phi$  in the Lie groupoid  $\Phi$  the curvature form of  $\tau$  and the curvature form of the underlying connection in  $A(\Phi)$  are differs on a constant

$$\Omega\left(\Psi_{2}^{\tau}\right) = \frac{1}{4} \cdot \Omega^{\tau}.$$

**Proof.** (a sketch) three steps of the proof:

-1) Of course, we need to prove the equality point by point, so we can look at this locally. Assume that  $\Phi$  is a trivial Lie groupoid  $\Phi = M \times G \times M$ with the source  $pr_3$  and the target  $pr_1$  and the partial multiplication

$$(z, a, y) \cdot (y, b, x) = (z, ab, x)$$

Then  $A(\Phi)_{|_{\mathbf{x}}} = T_{\mathbf{x}}M \times \mathfrak{g}$  where  $\mathfrak{g}$  is the **right** ! Lie algebra of G,

$$A\left(\Phi\right) = TM \times \mathfrak{g}.$$

The bracket in  $\mathfrak g$  will be denoted by  $[v,w]^R.$  The linear direct connection  $\tau$  is given by

$$\begin{aligned} \tau : (M \times M)_{|U} &\to M \times G \times M \\ \tau \left( x, y \right) = \left( x, \hat{\tau} \left( x, y \right), y \right), \\ \hat{\tau} : \left( M \times M \right)_{|U} &\to G, \quad \hat{\tau} \left( x, x \right) = e. \end{aligned}$$

Therefore

$$\frac{\partial \hat{\tau}}{\partial x^i|_{(x,x)}} + \frac{\partial \hat{\tau}}{\partial x^{m+i}|_{(x,x)}} = 0$$

and

$$\Psi_{2}^{\tau}(x_{0}, x_{1}, x_{2}) = (x_{0}, \ \hat{\tau}(x_{0}, x_{1}) \cdot \hat{\tau}(x_{1}, x_{2}) \cdot \hat{\tau}(x_{2}, x_{0}), \ x_{0}).$$

The induced linear connection  $\nabla^\tau:TM\to TM\times\mathfrak{g}$  is equal to

$$\nabla_{|y}^{\tau}(v) = \nabla_{v}^{\tau} = \tau\left(\cdot, y\right)_{*y}(v) = \left(v, \hat{\tau}\left(\cdot, y\right)_{*y}(v)\right).$$

We calculate the curvature tensor  $\Omega^{\tau}$  of  $\nabla^{\tau}$ .

$$\Omega_{|x_0}^{\tau} = 2\sum_{i < j} \left( \frac{\partial^2 \hat{\tau}}{\partial x^j \partial x^{m+i}} - \frac{\partial^2 \hat{\tau}}{\partial x^i \partial x^{m+j}} + \left[ \frac{\partial \hat{\tau}}{\partial x^i}, \frac{\partial \hat{\tau}}{\partial x^j} \right]^R \right)_{|(x_0, x_0)} dx^i \wedge dx^j.$$

-ii) Now we calculate the form  $\Omega(\Psi_2)$ :

$$\Omega\left(\Psi_{2}\right)\left(x_{0}\right) = \frac{1}{2!} \sum_{i,j} \frac{\partial^{2}\Psi_{2}\left(x_{0}, x_{1}, x_{2}\right)}{\partial x_{1}^{i} \partial \tilde{x}_{2}^{j}} |_{x_{0}=x_{1}=x_{2}} dx^{i} \wedge dx^{j}$$

for  $\Psi_2(x_0, x_1, x_2) = \hat{\tau}(x_0, x_1) \cdot \hat{\tau}(x_1, x_2) \cdot \hat{\tau}(x_2, x_0)$ .

-a) Firstly we calculate it for the Lie group G of matrices,

 $G \subset GL(V) \subset \mathbb{R}^{(\dim V)^2}$ 

for some finitely dimensional vector space V, using the fact that the differential of the left and the right translations,  $l_g$  and  $r_g$ , are exactly respectively the left and the right multiplication by matrices,

$$(l_g)_{*h} = gh, \ (r_g)_{*h} = hg.$$

After calculations we obtain the equality

$$\Omega\left(\Psi_2\right) = \frac{1}{4} \cdot \Omega^{\tau}.$$

-b) For an arbitrary Lie group G we use the following "IMPORTANT LOCAL TRICK":

We need to consider only local Lie group structure near the unit. But every Lie algebra is isomorphic to a Lie algebra of matrices (because there exists a faithful representation in some finitely dimensional vector space V) and a Lie algebra of matrices is a Lie algebra of a Lie subgroup of the Lie group GL(V). Therefore the above result concerning  $G \subset GL(V)$  is valid for arbitrary Lie group G!

#### 3.5 Characteristic classes

The last theorem gives that we can extract the Chern-Weil homomorphism of  $\Phi$  via any local direct connection  $\tau$  on the level of differential forms. The Chern-Weil homomorphism of  $\Phi$  is really the Chern-Weil homomorphism of the Lie algebroid  $A(\Phi)$  of  $\Phi$ .

We recall the construction of the Chern-Weil homomorphism for Lie algebroids

• Jan Kubarski, *The Chern-Weil homomorphism of regular Lie algebroids*, UNIVERSITE CLAUDE BERNARD – LYON 1, Publications du Départment de Mathématiques, nouvelle série, 1991, 1-70.

Consider a transitive Lie algebroid A with the Atiyah sequence

$$0 \to \boldsymbol{g} \to A \to TM \to 0$$

with the adjoint bundle of Lie algebras g.

The Chern-Weil homomorphism for transitive Lie algebroid A is defined as follows:

$$h_{A}: \bigoplus^{k \ge 0} \left( \operatorname{Sec} \bigvee^{k} \boldsymbol{g}^{*} \right)_{I^{0}} \longrightarrow \mathbf{H}_{dR} \left( M \right)$$
$$\Gamma \longmapsto \left[ \frac{1}{k!} \langle \Gamma, \Omega \lor \dots \lor \Omega \rangle \right]$$

where  $\Omega \in \Omega^2_E(M; \boldsymbol{g})$  is the curvature tensor of any connection in A, whereas  $\left(\operatorname{Sec} \bigvee^k \boldsymbol{g}^*\right)_{I^0}$  is the space of invariant cross-sections of  $\bigvee^k \boldsymbol{g}^*$  with respect to the adjoint representation of A on  $\bigvee^k \boldsymbol{g}^*$ , i.e.  $\Gamma \in \left(\operatorname{Sec} \bigvee^k \boldsymbol{g}^*\right)_{I^0}$  if and only if

$$\forall_{\xi \in \operatorname{Sec} A} \forall_{\sigma_1, \dots, \sigma_k \in \operatorname{Sec} \boldsymbol{g}} \left( (\gamma \circ \xi) \langle \Gamma, \sigma_1 \lor \dots \lor \sigma_k \rangle = \sum_{i=1}^k \langle \Gamma, \sigma_1 \lor \dots \lor \llbracket \xi, \sigma_i \rrbracket \lor \dots \lor \sigma_k \rangle \right)$$

The nontriviality of  $h_A$  means, of course, that in A there is no flat connection.

We explain also that  $\Omega \lor ... \lor \Omega$  is the usual skew multiplication of differential forms with values multiplying symmetrically

$$\Omega \lor \dots \lor \Omega\left(x; v_1, \dots, v_{2k}\right) = \sum_{\sigma \in \Sigma^{2k}} sgn \, \sigma \cdot \Omega\left(x; v_{\sigma_1}, v_{\sigma_2}\right) \lor \dots \lor \Omega\left(x; v_{\sigma_{2k-1}}, v_{gs_{2k}}\right) \in \bigvee^k \boldsymbol{g}_x.$$

For example: for the Lie algebroid A(P) of a principal fibre bundle P(P) is assumed to be connected) and equivalently for Lie algebroid of the Ehresmann Lie groupoid  $\Phi = PP^{-1}$ , there is a natural isomorphism of algebras  $\nu$  such that the diagram commutes

which means that the Chern-Weil homomorphism of a Lie algebroid is some generalization of this notion known on the ground of principal bundles. On the other hand, this also means that the Chern-Weil homomorphism of a principal bundle is a characteristic feature of its Lie algebroid (for connected principal bundles).

In addition, we must point out two things:

- 1) A Lie algebroid is in some sense a simpler structure than a principal bundle. Namely, nonisomorphic principal bundles can possess isomorphic Lie algebroids. For example, there exists a nontrivial principal bundle for which the Lie algebroid is trivial (the nontrivial Spin (3)-structure of the trivial principal bundle  $\mathbb{R}P(5) \times SO(3)$ ).
- 2) There exist other sources of Lie algebroids than principal bundles, for example, transversally complete foliations, nonclosed Lie subgroups, Poisson manifolds and other.

**Example 19** Let  $\Phi = GL(E)$  be a Lie groupoid of all linear fibre isomorphisms. Then the Atiyah sequence of  $A(\Phi)$  is

$$0 \to End(E) \to A(\Phi) \to TM \to 0.$$

Consider the Chern-Weil homomorphism

$$h: \bigoplus^{k\geq 0} \left(\operatorname{Sec} \bigvee^{k} End\left(E\right)^{*}\right)_{I^{0}} \to H_{dR}\left(M\right), \Gamma \longmapsto \left[\frac{1}{k!} \langle \Gamma, \Omega \lor \dots \lor \Omega \rangle\right]$$

where  $\Omega \in \Omega^2(M; End(E))$  is the curvature tensor of any connection in  $\Phi$ .

**Pontryagin classes.** Take the **invariant** cross section  $C_k \in \Gamma\left(\operatorname{Sec} \bigvee^k End\left(E\right)^*\right)$  by

$$C_{k|x} = tr\left(\varphi_1 \Box ... \Box \varphi_k\right)$$

where for  $\varphi_i \in End(E_{|x})$  the linear mapping  $\varphi_1 \Box ... \Box \varphi_k : \bigwedge^k E_{|x} \to \bigwedge^k E_{|x}$  is defined [Greub-Halperin-Vanstone] by

$$\varphi_{1}\Box\ldots\Box\varphi_{k}\left(v_{1}\wedge\ldots\wedge v_{k}\right)=\sum_{\sigma\in\Sigma^{k}}sgn\sigma\cdot\varphi_{1}\left(v_{\sigma_{1}}\right)\wedge\ldots\wedge\varphi_{k}\left(v_{\sigma_{k}}\right).$$

Then the Ponryagin class is equal to

$$p_k(E) = p_k(\Phi) = h(C_{2k}) = \frac{1}{(2k)!} \left[ \langle C_{2k}, \underbrace{\Omega \lor \dots \lor \Omega}_{2k \text{ times}} \rangle \right]_{dR}.$$

The class  $p_k(E)$  is represented by the differential form

$$c \cdot tr (\Omega \Box ... \Box \Omega)$$
.

According to the notation of Greub-Halperin-Vanstone, the forms  $\Omega \lor ... \lor \Omega$  and  $\Omega \Box ... \Box \Omega$  are the usual skew multiplication of differential forms for which the values are multiplicated by the suitable mappings

$$\forall : End(E) \times ... \times End(E) \rightarrow \bigvee^{2k} End(E) .$$
$$\Box : End(E) \times ... \times End(E) \rightarrow End\left(\bigwedge^{2k} E\right).$$

**Trace classes.** Take the **invariant** cross section  $Tr_k \in \Gamma\left(\operatorname{Sec} \bigvee^k End\left(E\right)^*\right)$  by

$$Tr_{k}\left(\varphi_{1},...,\varphi_{k}\right)=\sum_{\sigma\in\Sigma^{k}}tr\left(\varphi_{\sigma_{1}}\circ...\circ\varphi_{\sigma_{k}}\right).$$

Then the trace class is equal to

$$tr_{k}(E) = tr_{k}(\Phi) = h(Tr_{2k}) = \frac{1}{(2k)!} \left[ \langle Tr_{2k}, \underbrace{\Omega \lor \dots \lor \Omega}_{2k \text{ times}} \rangle \right]_{dR}.$$

The class  $tr_k(E)$  is represented by the differential form

 $c \cdot tr \left(\Omega \circ \dots \circ \Omega\right)$ 

(the values of the skew multiplication of  $\Omega \circ ... \circ \Omega$  are multiplied by the composing of the linear mapping.

**Pffafian class for oriented** 2k-dimensional vector bundle E. Take the invariant cross-section  $pf \in \Gamma\left(\operatorname{Sec}\bigvee^k Sk\left(E\right)^*\right)$ 

$$pf^{F}\left(\varphi_{1},...,\varphi_{k}\right)=\left\langle e,\beta^{-1}\left(\varphi_{1}\right)\wedge...\wedge\beta^{-1}\left(\varphi_{k}\right)\right\rangle$$

where  $\beta : \bigwedge^2 (F) \xrightarrow{\cong} Sk_F$ ,  $\beta (x \wedge y)(z) = \langle x, z \rangle y - \langle y, z \rangle x$  and  $e \in \bigwedge^k F$  determine the orientation and  $|\langle e, e \rangle| = 1$ . Then the Atiyah sequence of A(IsoE) is

 $0 \to Sk(E) \to A(IsoE) \to TM \to 0.$ 

and the Pfaffian class is equal to  $i^k \cdot h(pf)$  and it is represented by

$$c \cdot \langle \Delta, (\beta^{-1}\Omega) \wedge ... \wedge (\beta^{-1}\Omega) \rangle.$$

Take a local direct connection  $\tau$  in  $\Phi$  and consider once again the curvature form  $\Omega(\Psi_2^{\tau}) \in \Omega^2(M; \boldsymbol{g})$ ,

$$\Omega_{|x_0}^{\tau} = 2\sum_{i < j} \left( \frac{\partial^2 \hat{\tau}}{\partial x^j \partial x^{m+i}} - \frac{\partial^2 \hat{\tau}}{\partial x^i \partial x^{m+j}} + \left[ \frac{\partial \hat{\tau}}{\partial x^i}, \frac{\partial \hat{\tau}}{\partial x^j} \right]^R \right)_{|(x_0, x_0)} dx^i \wedge dx^j.$$

We known that  $\tau$  induces a usual connection  $\nabla^{\tau}$  in  $A(\Phi)$  and that the curvature of it is a constant time the form  $\Omega(\Psi_2^{\tau})$ ,

$$\Omega^{\tau} = 4 \cdot \Omega \left( \Psi_2^{\tau} \right).$$

In conclusion, the Chern-Weil homomorphism of  $\Phi$  (i.e. of the  $A(\Phi)$ ) can be extracted via  $\tau$  on the **level of differential forms** by

$$\langle \Gamma, \Omega \lor \ldots \lor \Omega \rangle = 4^k \langle \Gamma, \Omega \left( \Psi_2^{\tau} \right) \lor \ldots \lor \Omega \left( \Psi_2^{\tau} \right) \rangle.$$

Problem 20 How can we express the form

$$\Omega\left(\Psi_{2}^{\tau}\right) \lor \ldots \lor \Omega\left(\Psi_{2}^{\tau}\right) \in \Omega^{2k}\left(M; \bigvee \boldsymbol{g}\right)$$

with the help of  $\Omega(\Psi_{2k}^{\tau})$ ? and the form  $\langle \Gamma, \Omega(\Psi_{2}^{\tau}) \vee ... \vee \Omega(\Psi_{2}^{\tau}) \rangle$  for an invariant cross-section  $\Gamma \in \left( \operatorname{Sec} \bigvee^{k} \boldsymbol{g}^{*} \right)_{I^{0}}$  with the help  $\Gamma$  and  $\Omega(\Psi_{2k}^{\tau})$ ?

We know that for  $\Phi = GL(E)$  we have the adjoint Lie algebra bundle  $\boldsymbol{g}$  is equal to the vector bundle of linear homomorphisms Aut(E). Therefore the Atiyah sequence of  $A(\Phi)$  equals

$$0 \to Aut(E) \to A(\Phi) \to TM \to 0.$$

Using the composition of linear homomorphisms

$$\circ: Aut\left(E\right) \times \dots \times Aut\left(E\right) \to Aut\left(E\right)$$

we can obtain the following theorem.

**Theorem 21** The equality holds

$$\Omega\left(\Psi_{2k}^{\tau}\right) = c \cdot \Omega\left(\Psi_{2}^{\tau}\right) \circ \dots \circ \Omega\left(\Psi_{2}^{\tau}\right) \quad (k \text{ times})$$

or equivalently

$$\Omega\left(\Psi_{2k}^{\tau}\right) = c_1 \cdot \Omega^{\tau} \circ \dots \circ \Omega^{\tau} \quad (k \ times).$$

THE END

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Remark concerning the above Problem (20): For an arbitrary Lie groupoid  $\Phi$  we consider

$$\Psi_{2k}^{\tau}\left(x_{0}, x_{1}, ..., x_{2k}\right) = \tau\left(x_{0}, x_{1}\right) \cdot \tau\left(x_{1}, x_{2}\right) \cdot ... \cdot \tau\left(x_{2k-1}, x_{2k}\right) \cdot \tau\left(x_{2k}, x_{0}\right) \in \Phi_{x_{0}}^{x_{0}}.$$

We define additionally

$$\begin{split} \tilde{\Psi}_{2k}^{\tau} \left( x_0, x_1, \dots, x_{2k} \right) \\ &= \tau \left( x_0, x_1 \right) \cdot \tau \left( x_1, x_2 \right) \cdot \left[ \tau \left( x_2, x_0 \right) \cdot \tau \left( x_0, x_3 \right) \right] \cdot \tau \left( x_3, x_4 \right) \cdot \left[ \tau \left( x_4, x_0 \right) \cdot \tau \left( x_0, x_5 \right) \right] \cdot \dots \\ &\dots \cdot \left[ \tau \left( x_{2k-2}, x_{2k} \right) \cdot \tau \left( x_0, x_{2k-1} \right) \right] \cdot \tau \left( x_{2k-1}, x_{2k} \right) \cdot \tau \left( x_{2k}, x_0 \right) \\ &= \left[ \tau \left( x_0, x_1 \right) \cdot \tau \left( x_1, x_2 \right) \cdot \tau \left( x_2, x_0 \right) \right] \cdot \left[ \tau \left( x_0, x_3 \right) \cdot \tau \left( x_3, x_4 \right) \cdot \tau \left( x_4, x_0 \right) \right] \cdot \dots \\ &\dots \cdot \left[ \tau \left( x_0, x_{2k-1} \right) \cdot \tau \left( x_{2k-1}, x_{2k} \right) \cdot \tau \left( x_{2k}, x_0 \right) \right] \end{split}$$

i.e.

$$\tilde{\Psi}_{2k}^{\tau}(x_0, x_1, \dots, x_{2k}) = \Psi_2^{\tau}(x_0, x_1, x_2) \cdot \Psi_2^{\tau}(x_0, x_3, x_4) \cdot \dots \cdot \Psi_2^{\tau}(x_0, x_{2k-1}, x_{2k})$$

Clearly  $\Psi_{2k}^{\tau} \neq \tilde{\Psi}_{2k}^{\tau}$  in general, but (I think) the equality holds infinitesimally

$$\Omega\left(\Psi_{2k}^{\tau}\right)\left(x\right) = \Omega\left(\tilde{\Psi}_{2k}^{\tau}\right)\left(x\right). \tag{3}$$

Next, using (3), we can try to calculate  $\Omega(\Psi_{2k}^{\tau})(x)$  via  $\Omega(\Psi_{2}^{\tau})(x)$ .